



Search for dibaryon resonances in the reactions $dd \rightarrow dd\pi\pi$ and $pd \rightarrow pd\pi\pi$

Yu. Uzikov¹⁾

in collaboration with A. Temerbayev²⁾ and N. Tursunbayev²⁾

¹⁾ *V.P. Dzhelepov Laboratory of Nuclear Problems, JINR*

uzikov@jinr.ru

²⁾ *L.N. Gumilyov Eurasian National University, Astana, Kazakhstan*

Baldin ISHEPP XXVI, 15-20 September 2025, Dubna

CONTENT

- WASA dibaryon $d^*(2380)$ in the reactions $pn \rightarrow d\pi^0\pi^0$ and $pn \rightarrow pn$. Experimental data and theoretical models.
- ANKE@COSY data on $pd \rightarrow pd\pi\pi$ and t-channel excitation of $d^*(2380)$.
- Indication to isoscalar resonances in $\gamma + d \rightarrow d\pi^0\pi^0$ (ELPH and BGOOD data).
- About a possibility to search for isoscalar dibaryons in $dd \rightarrow dd \pi\pi$ and $pd \rightarrow pd\pi\pi$ reactions in dd-collision at SPD NICA .
- Summary

Table I. $Y = 2$ states with zero strangeness predicted by the 490 multiplet.

Particle	T	J	SU(3) multiplet	Comment	Predicted mass
D_{01}	0	1	<u>10*</u>	Deuteron	A
D_{10}	1	0	<u>27</u>	Deuteron singlet state	A
D_{12}	1	2	<u>27</u>	S -wave $N-N^*$ resonance	$A + 6B$
D_{21}	2	1	<u>35</u>	Charge-3 resonance	$A + 6B$
D_{03}	0	3	<u>10*</u>	S -wave N^*-N^* resonance	$A + 10B$
D_{30}	3	0	<u>28</u>	Charge-4 resonance	$A + 10B$

D_{TJ}

$M \sim A + B[T(T+1) + J(J+1) - 2]$,
 $B = 47$ MeV with D_{12} from $pp \leftrightarrow \pi^+ d$

$M(D_{03}) = 2350$ MeV

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ДИБАРИОННЫЕ РЕЗОНАНСЫ

M. M. Макаров

Reference to PRL 42,
№20 (1979) 1321;
 $\gamma d \rightarrow np$

Angular Dependence of Proton Polarization in the Reaction $\gamma d \rightarrow pn$ and
a Partial-Wave Analysis of Possible Dibaryon Resonances

H. Ikeda,^(a) I. Arai, H. Fujii, T. Fujii, H. Iwasaki, N. Kajiura,^(a) T. Kamae
K. Nakamura, T. Sumiyoshi, and H. Takeda^(b)
Department of Physics, University of Tokyo, Tokyo 113, Japan

and

K. Ogawa
National Laboratory for High Energy Physics, Tsukuba, Ibaraki 300-32, Japan

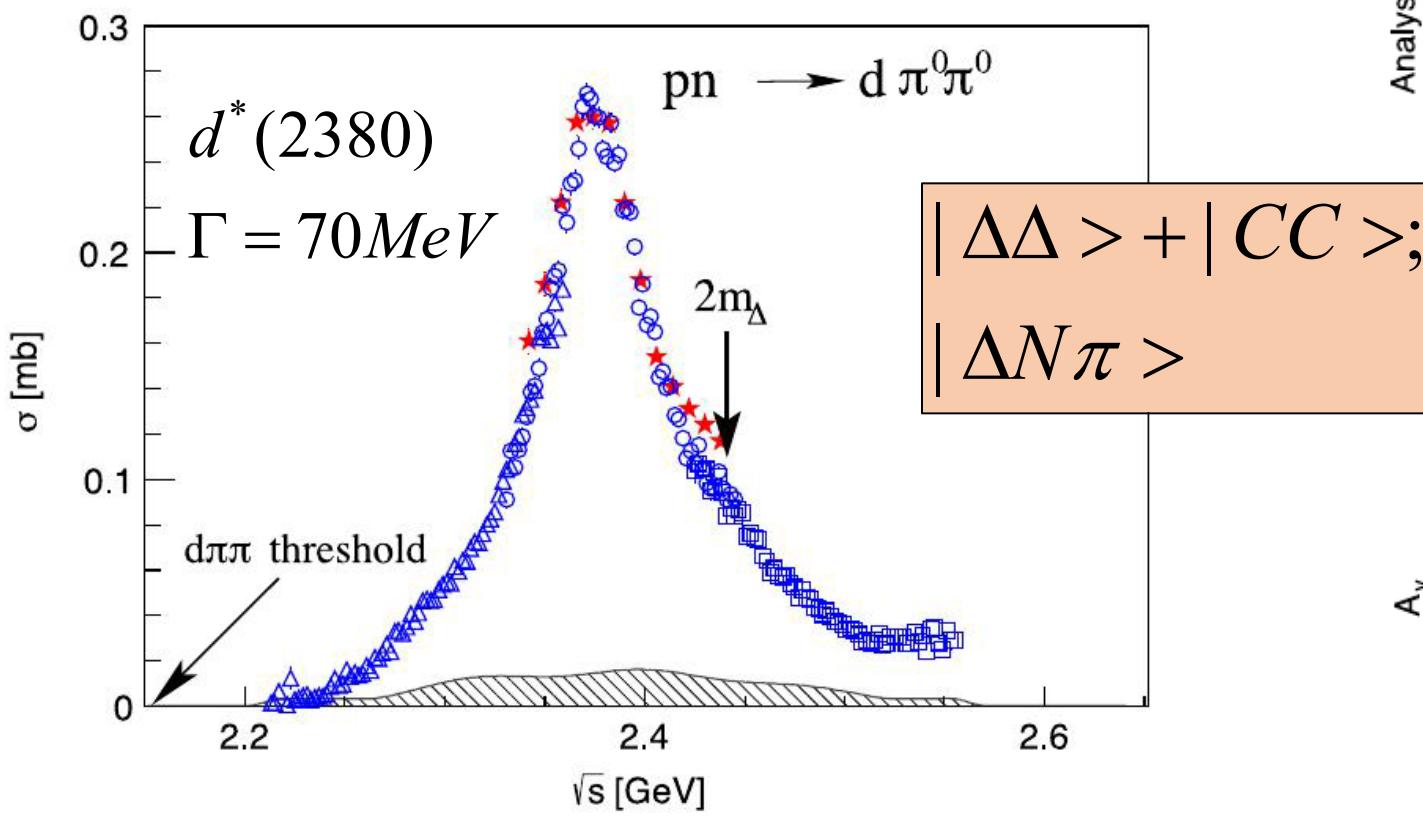
and

M. Kanazawa
Institute for Nuclear Study, University of Tokyo, Tokyo 188, Japan
(Received 22 March 1979)

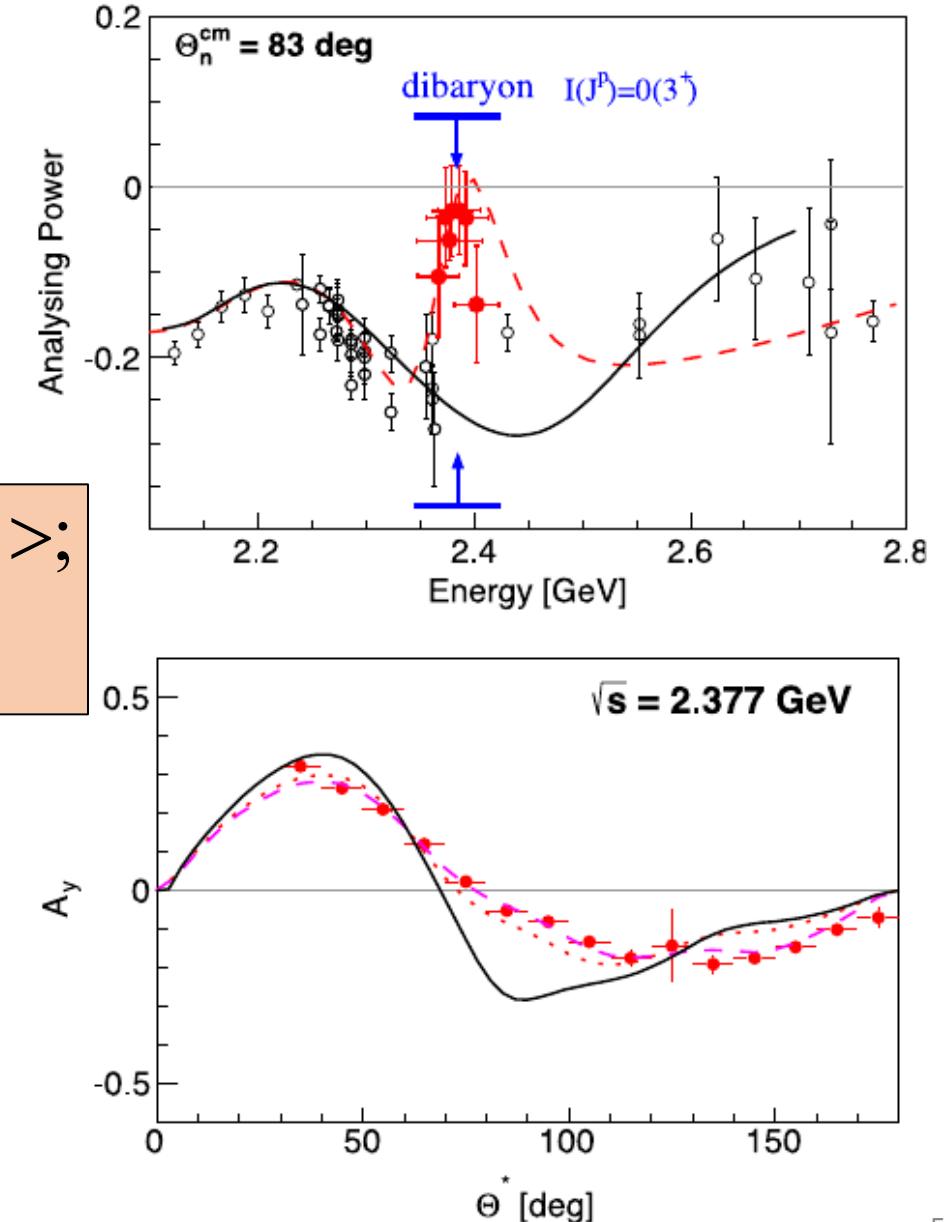
The angular dependence of proton polarization in $\gamma d \rightarrow pn$ has been measured at photon energies between 400 and 650 MeV. The polarization and differential-cross-section data are consistently explained by introducing a dibaryon resonance $I(J^P) = 0(3^+)$ or $0(1^+)$ at ≈ 2360 MeV.

DIBARYON RESONANCES

H. Clement / Progress in Particle and Nuclear Physics 93 (2017) 195–242



18.09.2025



5

H. Clement, Prog. Part. Nucl. Phys. 93 (2017) 195-242

H. Clement, T. Skorodko. Chin. Phys. C 45 , № 2 (2021) 022001

Y. Dong et al. “d*(2380) in chiral constituent quark model”

Prog. Part. Nucl. Phys. 131 (2023) 104045:

Compact hexaquark system, hidden color component dominates.

Eur. Phys. J. C (2023) 83:645

<https://doi.org/10.1140/epjc/s10052-023-11814-2>

$$d^* \sim -\sqrt{\frac{1}{3}} |\Delta\Delta\rangle + \sqrt{\frac{2}{3}} |CC\rangle$$

*See talk by Yubing Dong
on 16 September*

Regular Article - Theoretical Physics

Is $d^*(2380)$ a compact hexaquark state?

Manying Pan^{1,a}, Xinmei Zhu^{2,b}, Jialun Ping^{1,c}

However,

A.Gal, H. Garcilazo, Phys.Rev.Lett. 111 (2013) 172301; $d^*(2380)$ is the $\pi\Delta N$

See also talk by E. Doroshkevich, on 3 July, Nucleus-2025: NN*(1440), M=2315, Gamma=150 MeV

ANKE@COSY data on pd->pd $\pi\pi$ and t-channel excitation of d*(2380)

V.I. Komarov et al., Eur. Phys. Jour. A 54 (2018) 206

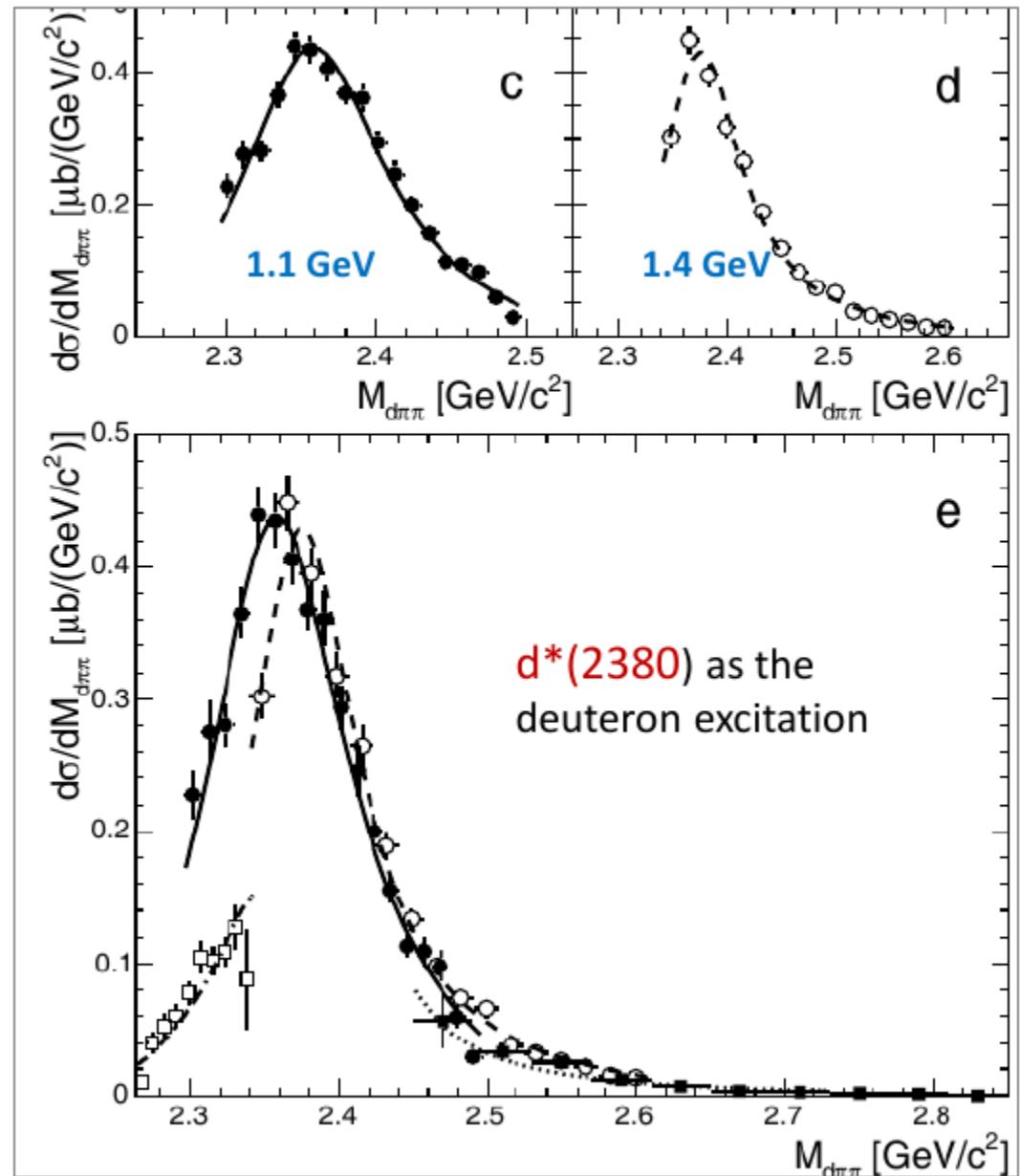
Eur. Phys. J. A (2018) 54: 206
DOI 10.1140/epja/i2018-12641-0

Regular Article – Experimental Physics

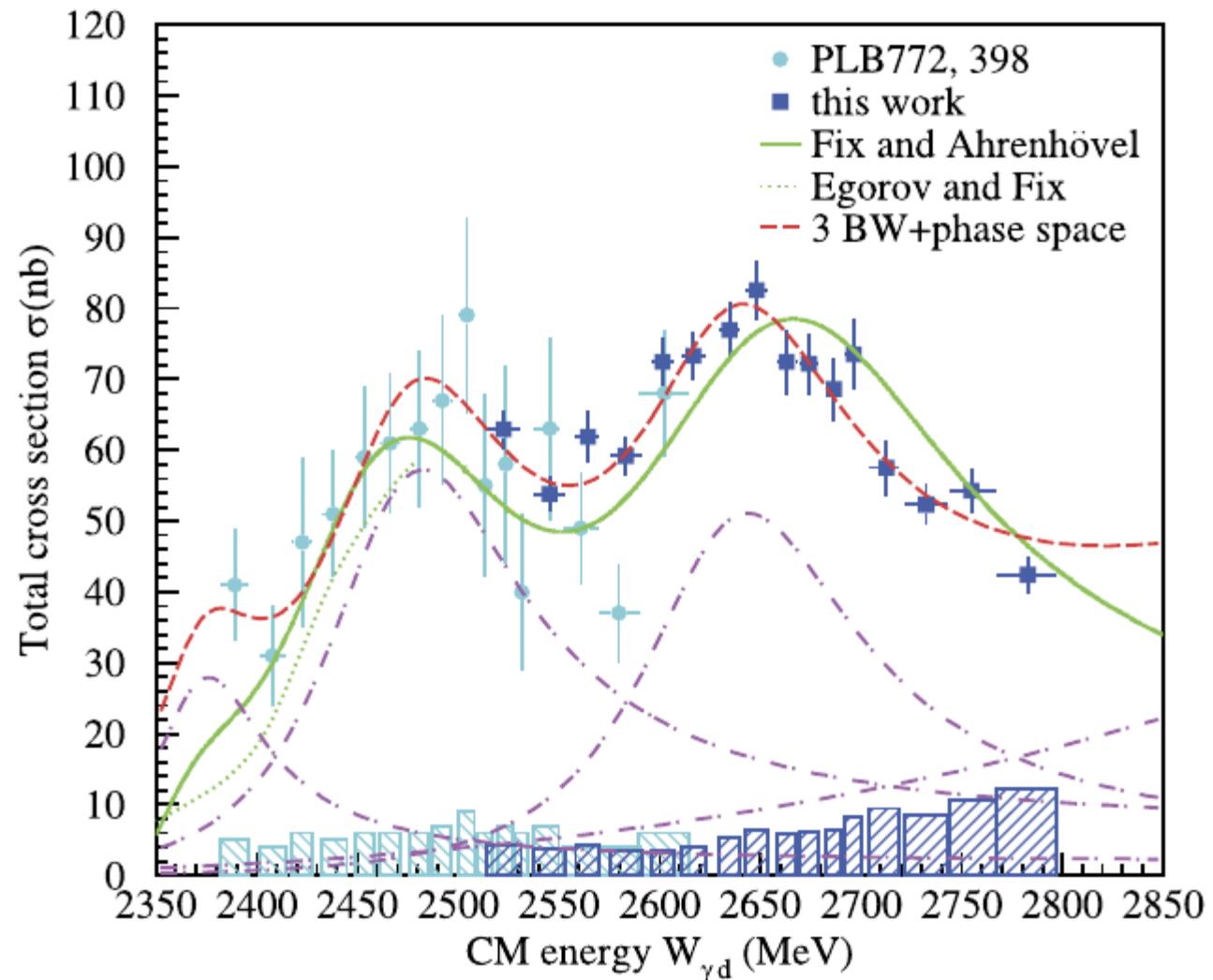
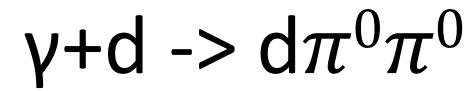
THE EUROPEAN
PHYSICAL JOURNAL A

Resonance-like coherent production of a pion pair in the reaction $pd \rightarrow pd\pi\pi$ in the GeV region

V.I. Komarov¹, D. Tsirkov^{1,a}, T. Azaryan¹, Z. Bagdasarian^{2,3}, B. Baimurzinova^{4,5}, S. Barsov⁶, S. Dymov^{1,2}, R. Gebel², M. Hartmann², A. Kacharava², A. Khoukaz⁷, A. Kulikov¹, A. Kunsafina^{1,4,5}, V. Kurbatov¹, Zh. Kurmanaliyev^{1,4,5}, B. Lorentz², G. Macharashvili³, D. Mchedlishvili^{2,3}, S. Merzliakov², S. Mikirtychians^{2,6}, M. Nioradze³, H. Ohm², F. Rathmann², V. Serdyuk², V. Shmakova¹, H. Ströher², S. Trusov^{2,8}, Yu. Uzikov^{1,9,10}, Yu. Valdau^{2,6}, and C. Wilkin¹¹



T. Ishikawa et al., (ELPH Coll) PLB 798 (2019) 413:
 $M_D = 2380, 2470, 2630$ MeV



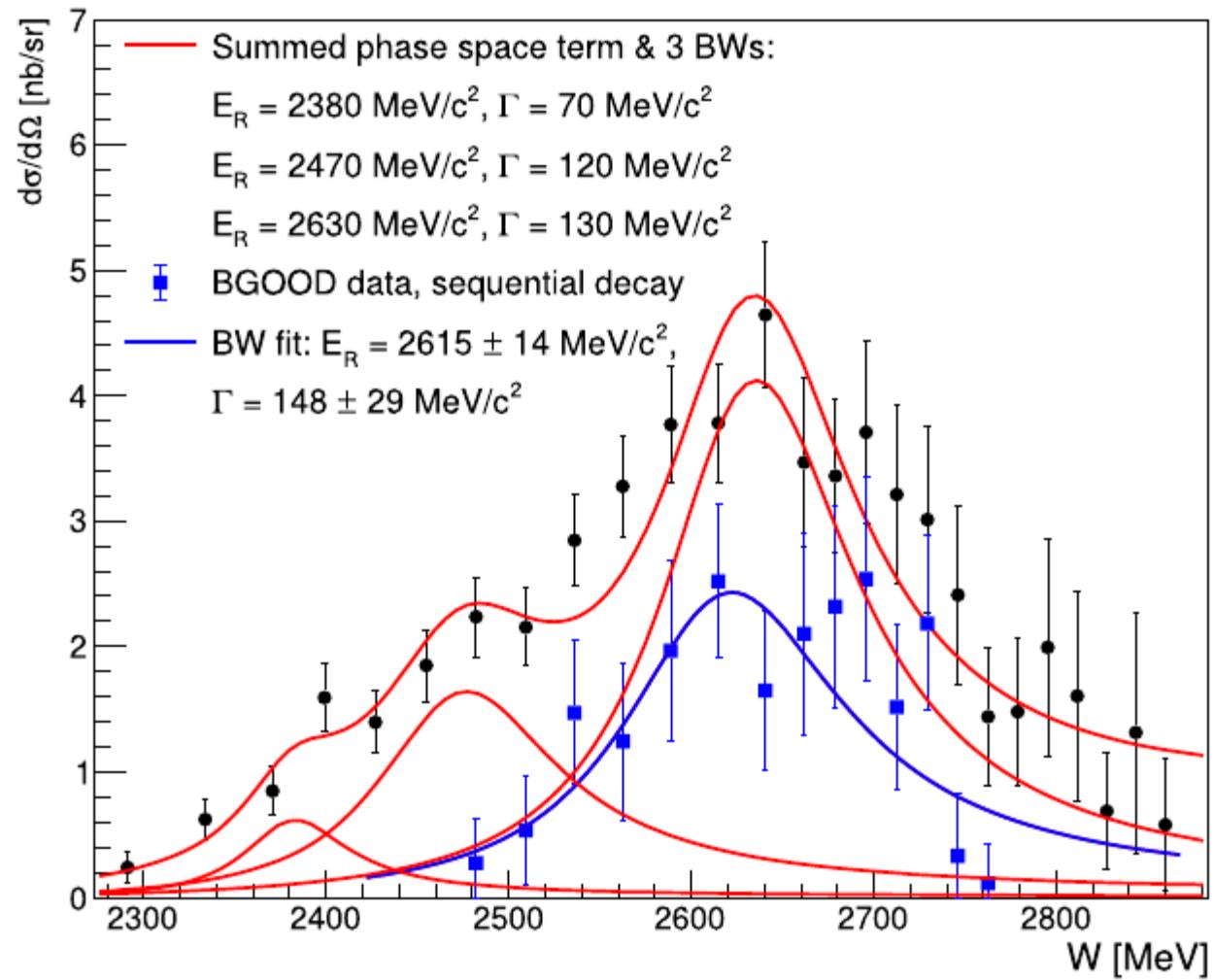


Fig. 7. $\gamma d \rightarrow \pi^0\pi^0d$ differential cross section for $\cos\theta_{CM}^d > 0.8$ (the same as in Fig. 5). A fit including three Breit-Wigner functions (BW) are shown as the red lines, with the fixed masses and widths labelled inset. The additional small centre-



$$\pi^0\pi^0 : I=0, 2$$

Isoscalar resonances: $I=0$

What is the origin of
the 2470 , 2630 MeV resonances?
Relation to the 2380 MeV?

Possible search at SPD in $dd - dd\pi\pi$
and $pd - pd\pi\pi$?
/B.F. Kostenko et al., weekly SPD meeting,
April 2025/

Isospin symmetry breaking in double-pion production in the region of $d^*(2380)$ and the scalar σ meson

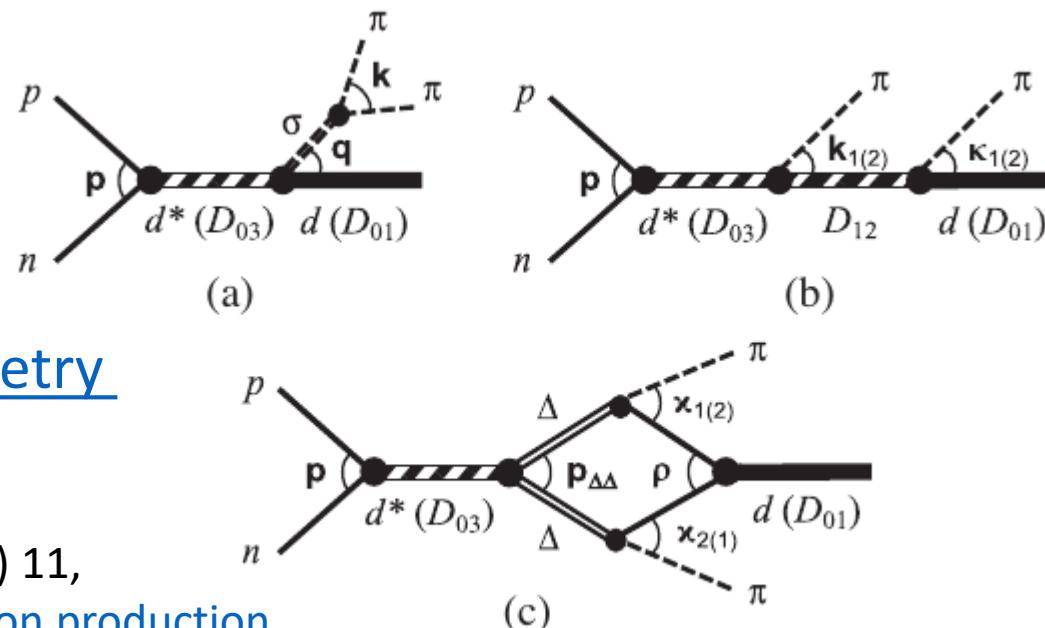
M. N. Platonova^{ID}[†] and V. I. Kukulin^{ID}^{*}

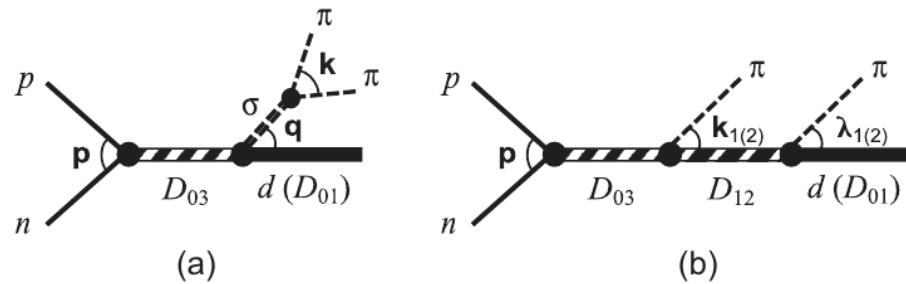
Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow 119991, Russia

M.P. Platonova, V.I. Kukulin,
PRC 87 (2013) 2, 025202

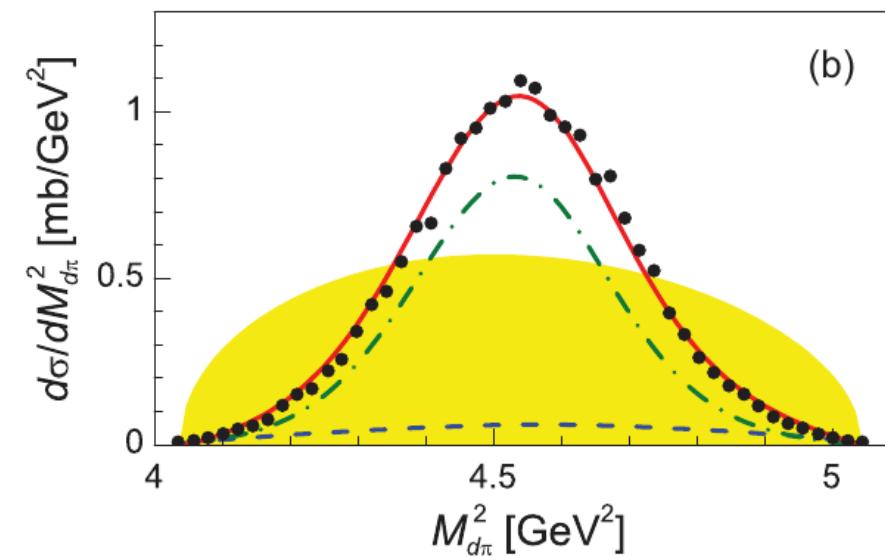
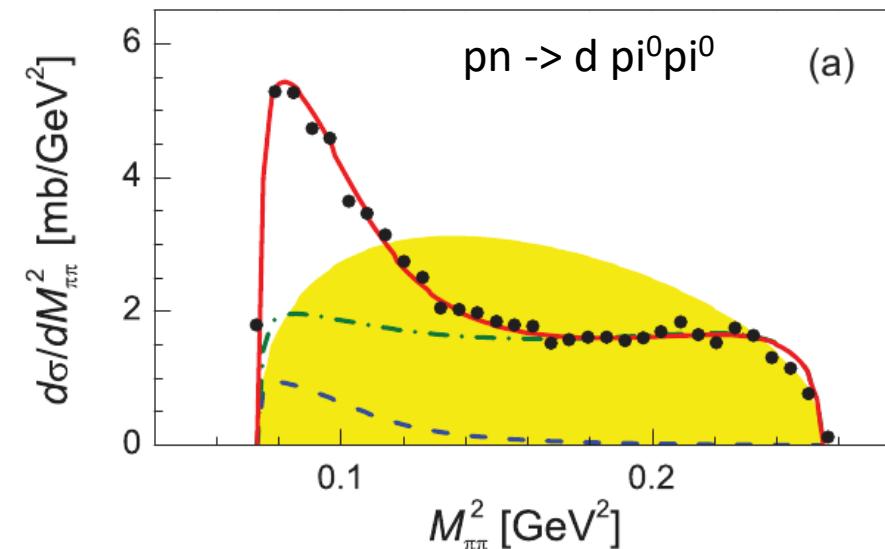
[ABC-effect as a signal of chiral symmetry
restoration in hadronic collisions](#)

M.P. Platonova, V.I. Kukulin, *Phys. Rev. D* 103 (2021) 11,
114025, [Isospin symmetry breaking in double-pion production
in the region of \$d^*\(2380\)\$ and the scalar \$\sigma\$ meson](#)

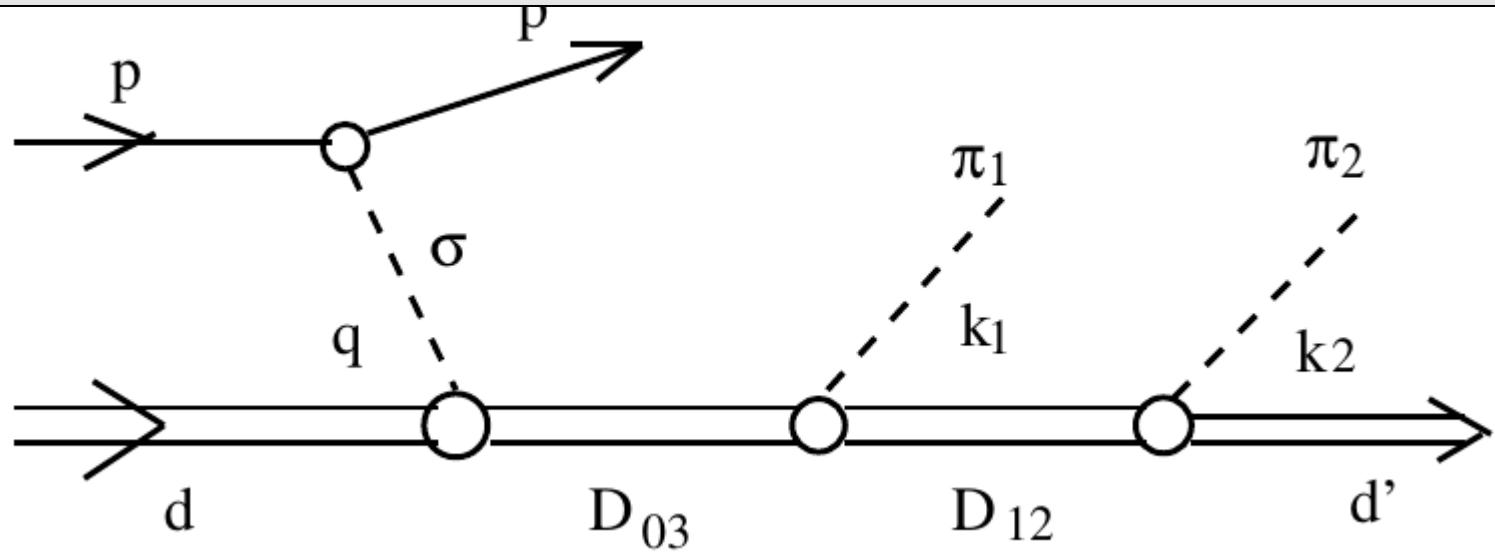




*M.N. Platonova, V. I. Kukulin, PRC 87 (2013)
250202*



Our Model of the reaction $pd \rightarrow pd\pi\pi$: PK + σ -exchange in t-channel



$$M_{\lambda_p \lambda_d}^{\lambda'_p \lambda'_d}(pd \rightarrow pd\pi\pi) = M_{\lambda_p}^{\lambda'_p}(p \rightarrow p'\sigma) \frac{1}{p_\sigma^2 - m_\sigma^2 + im_\sigma \Gamma_\sigma} M_{\lambda_d}^{\lambda'_d}(\sigma d \rightarrow d\pi\pi),$$

Yu. N. Uzikov, N. Tursunbayev, EPJ Web of Conf. **204** (2019) 08010

$$M_{\lambda_p \lambda_d}^{\lambda'_p \lambda'_d}(pd \rightarrow pd\pi\pi) = M_{\lambda_p}^{\lambda'_p}(p \rightarrow p'\sigma) \frac{1}{p_\sigma{}^2 - m_\sigma{}^2 + im_\sigma\Gamma_\sigma} M_{\lambda_d}^{\lambda'_d}(\sigma d \rightarrow d\pi\pi),$$

$$\begin{aligned} M_{\lambda_d}^{\lambda'_d}(\sigma d \rightarrow d\pi\pi) &= \sum_{\lambda_2, \lambda_3, \mu, m_1, m_2} \frac{F_{D_{03} \rightarrow d\sigma} F_{D_{03} \rightarrow D_{12}\pi_1}}{P_{D_{03}}^2 - M_{D_{03}}^2 + iM_{D_{03}}\Gamma_{D_{03}}} \frac{F_{D_{12} \rightarrow d\pi_2}}{P_{D_{12}}^2 - M_{D_{12}}^2 + iM_{D_{12}}\Gamma_{D_{12}}} \\ &\times (1\lambda_d 2\mu | 3\lambda_3) \mathcal{Y}_{2\mu}(\hat{\mathbf{q}}) (2\lambda_2 1m_1 | 3\lambda_3) \mathcal{Y}_{1m_1}(\hat{\mathbf{k}}_1) (1\lambda'_d 1m_2 | 2\lambda_2) \mathcal{Y}_{1m_2}(\hat{\mathbf{k}}_2) + (\pi_1 \leftrightarrow \pi_2), \end{aligned}$$

$$\frac{F_{D_{03} \rightarrow d\sigma}(q)}{M_{D_{03}}(q)} = \sqrt{\frac{8\pi\Gamma_{D_{03} \rightarrow d\sigma}^{(l=2)}(q)}{q^5}},$$

$$\Gamma_{D_{03} \rightarrow d\sigma}^{(l=2)}(q) = \Gamma_{D_{03} \rightarrow d\sigma}^{(l=2)}\left(\frac{q}{q_0}\right)^5 \left(\frac{q_0^2 + \lambda_{d\sigma}^2}{q^2 + \lambda_{d\sigma}^2}\right)^3,$$

$$\frac{F_{D_{03} \rightarrow D_{12}\pi_1}(k_1)}{M_{D_{12}\pi}(k_1)} = \sqrt{\frac{8\pi\Gamma_{D_{03} \rightarrow D_{12}\pi}^{(l=1)}(k_1)}{k_1^3}},$$

$$\Gamma_{D_{03} \rightarrow D_{12}\pi}^{(l=1)}(k_1) = \Gamma_{D_{03} \rightarrow D_{12}\pi}^{(l=1)}\left(\frac{k_1}{k_{10}}\right)^3 \left(\frac{k_{10}^2 + \lambda_{D_{12}\pi}^2}{k_1^2 + \lambda_{D_{12}\pi}^2}\right)^2$$

$$\frac{F_{D_{12} \rightarrow d\pi_2}(k_2)}{M_{d\pi_2}(k_2)} = \sqrt{\frac{8\pi\Gamma_{D_{12} \rightarrow d\pi}^{(l=1)}(k_2)}{k_2^3}},$$

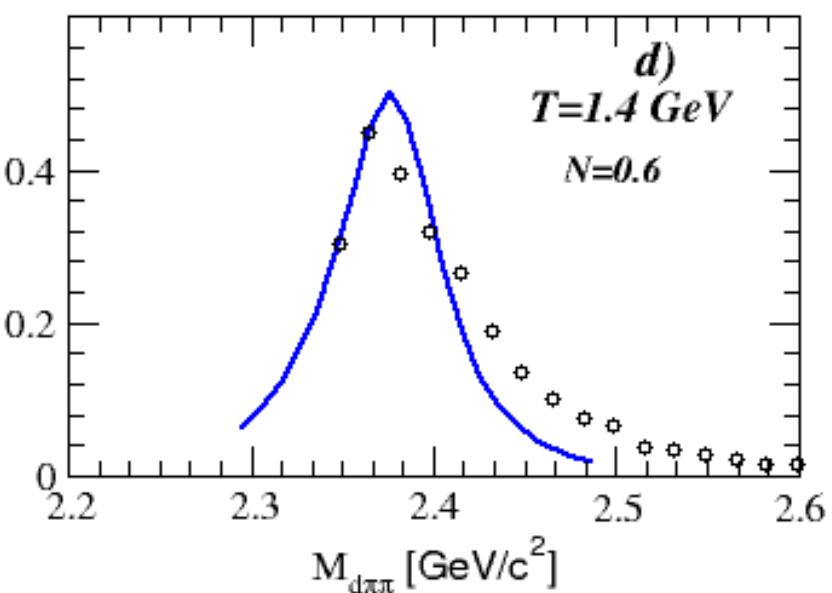
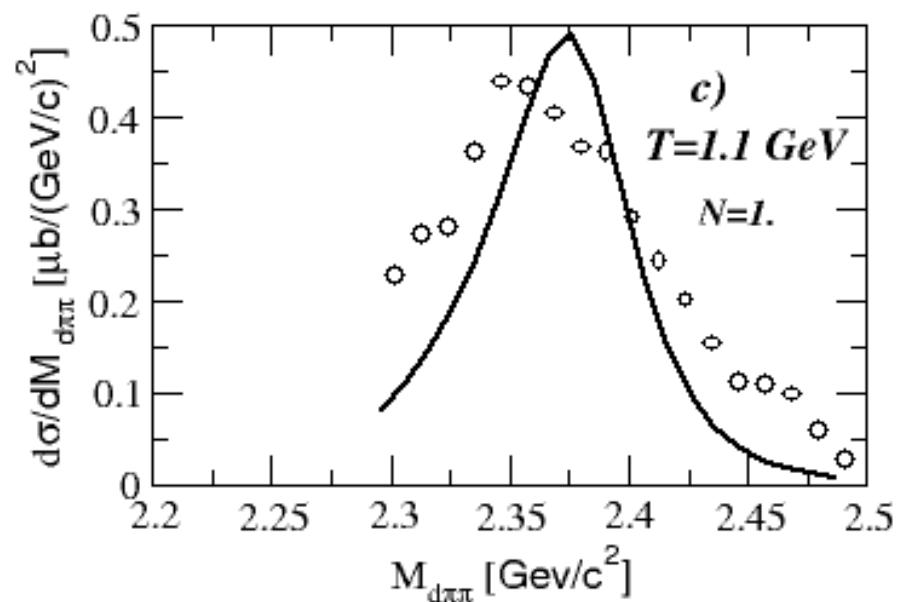
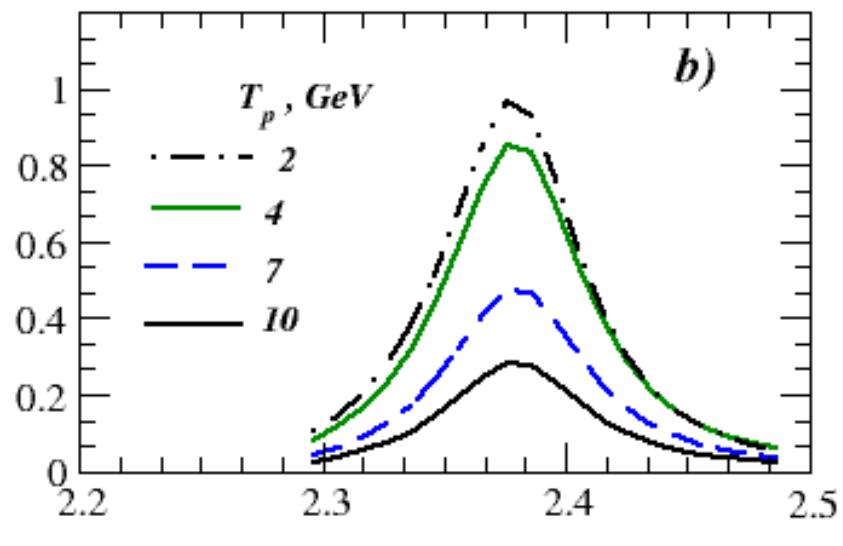
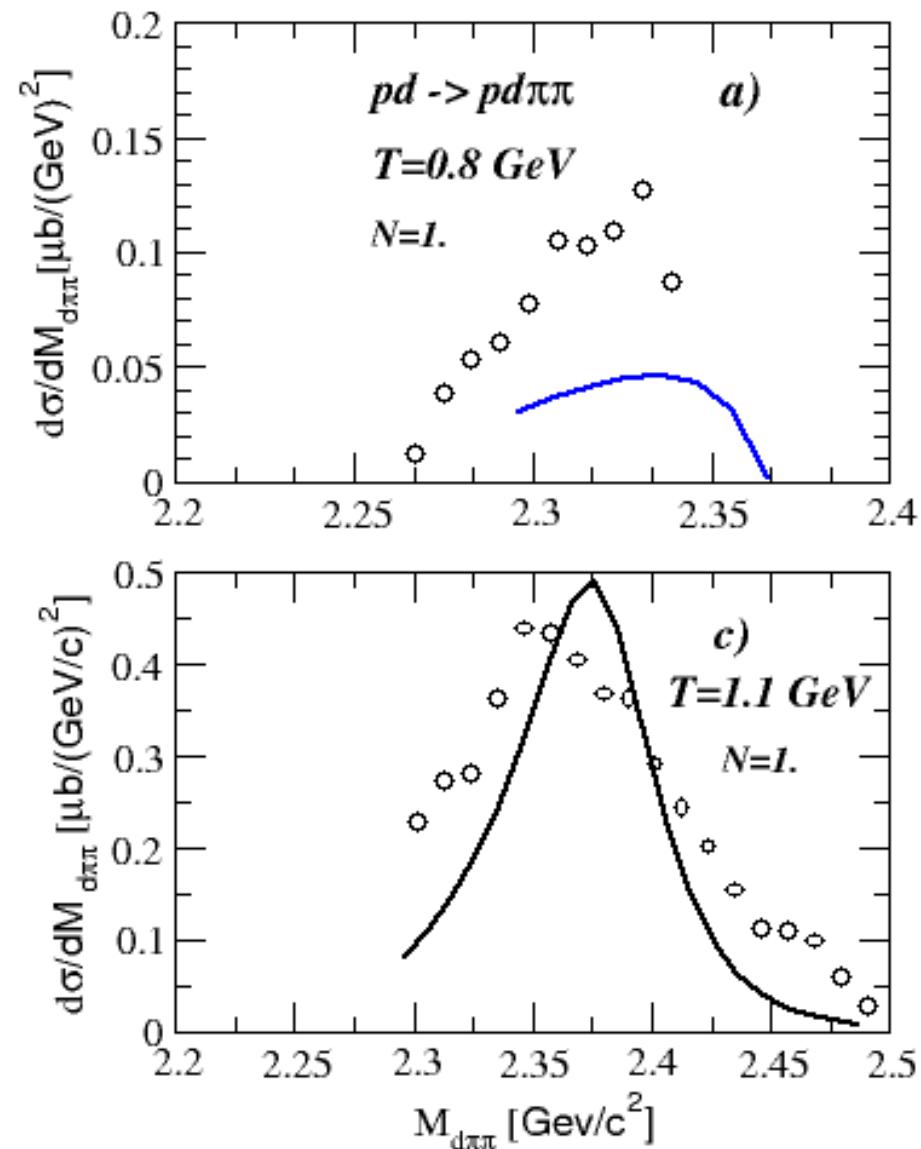
$$\Gamma_{D_{12} \rightarrow d\pi}^{(l=1)}(k_1) = \Gamma_{D_{12} \rightarrow d\pi}^{(l=1)}\left(\frac{k_2}{k_{20}}\right)^3 \left(\frac{k_{20}^2 + \lambda_{d\pi}^2}{k_2^2 + \lambda_{d\pi}^2}\right)^2.$$

$$\Gamma(p) = \Gamma(p_0) \left(\frac{p}{p_0}\right)^{2l+1} \left(\frac{p_0^2 + \Lambda^2}{p^2 + \Lambda^2}\right)^{l+1}$$

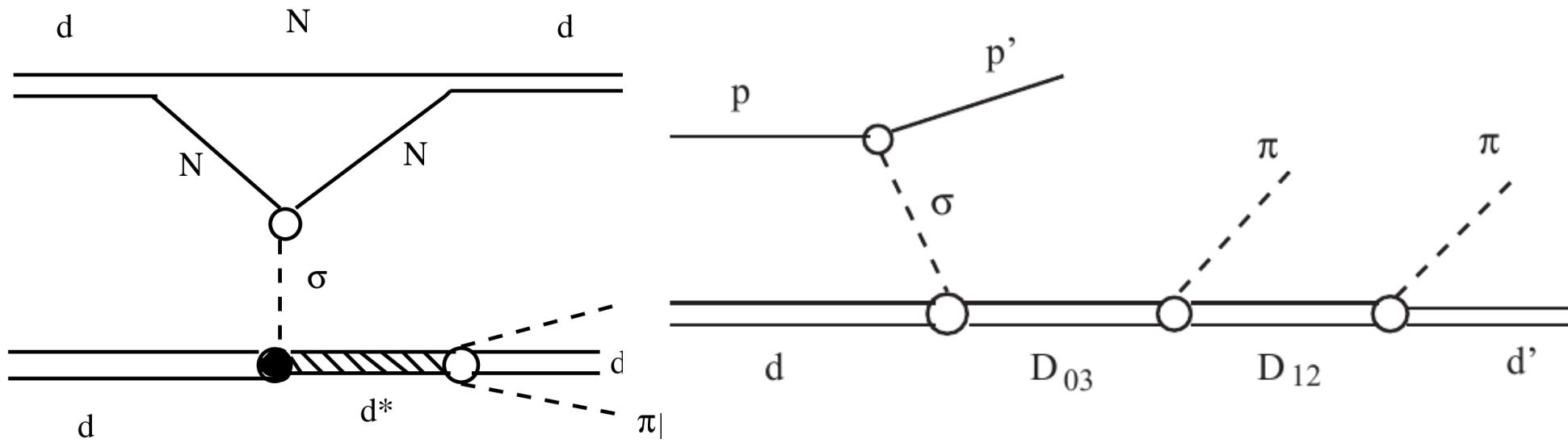
The following numbers were used in our calculations for the D_{03} and D_{12} resonances and the vertices parameters: $M_{D_{03}} = 2.380$ GeV, $\Gamma_{D_{03}} = 70$ MeV, $M_{D_{12}} = 2.15$ GeV, $\Gamma_{D_{12}} = 0.11$ GeV, $m_\sigma = 0.5$ GeV, $\Gamma_\sigma = 0.55$ GeV, $q_0 = 0.362$ Gev/c, $k_{10} = 0.177$ GeV/c, $k_{20} = 0.224$ GeV/c, $\lambda_{\pi D_{12}} = 0.12$ GeV. The values $\Gamma_{D_{12} \rightarrow d\pi}^{(l=1)} = 10$ MeV, $\lambda_{d\sigma} = 0.18$ GeV and $\lambda_{d\pi} = 0.25$ GeV were

$pd \rightarrow pd\pi\pi$: data - ANKE@COSY, theory - PK model + sigma in t-exchange

$\Theta^p_{cm} = 7^\circ - 13^\circ$



Model of the reaction $dd \rightarrow dd\pi\pi$ based on the $pd \rightarrow pd\pi\pi$ model

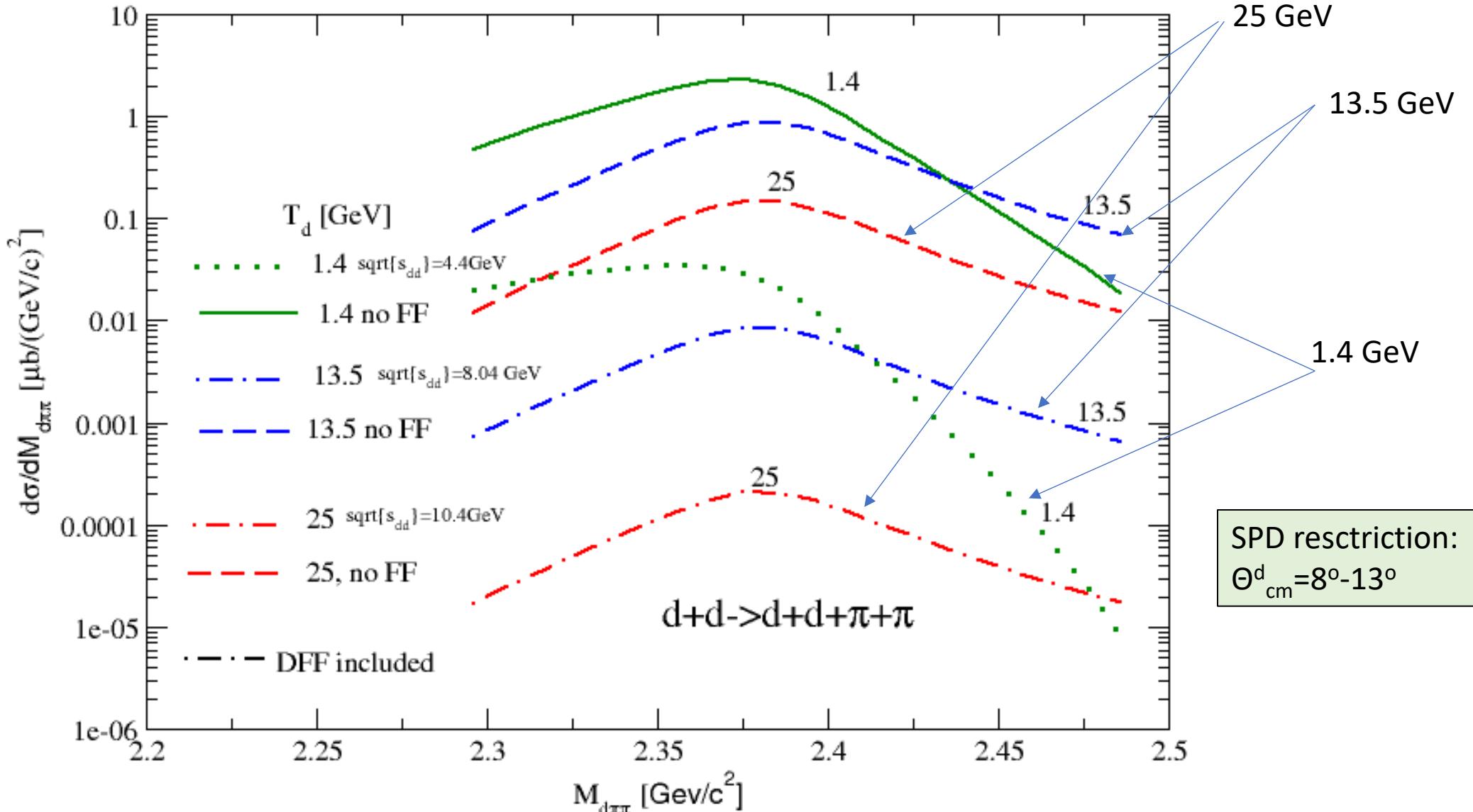


a)

b)

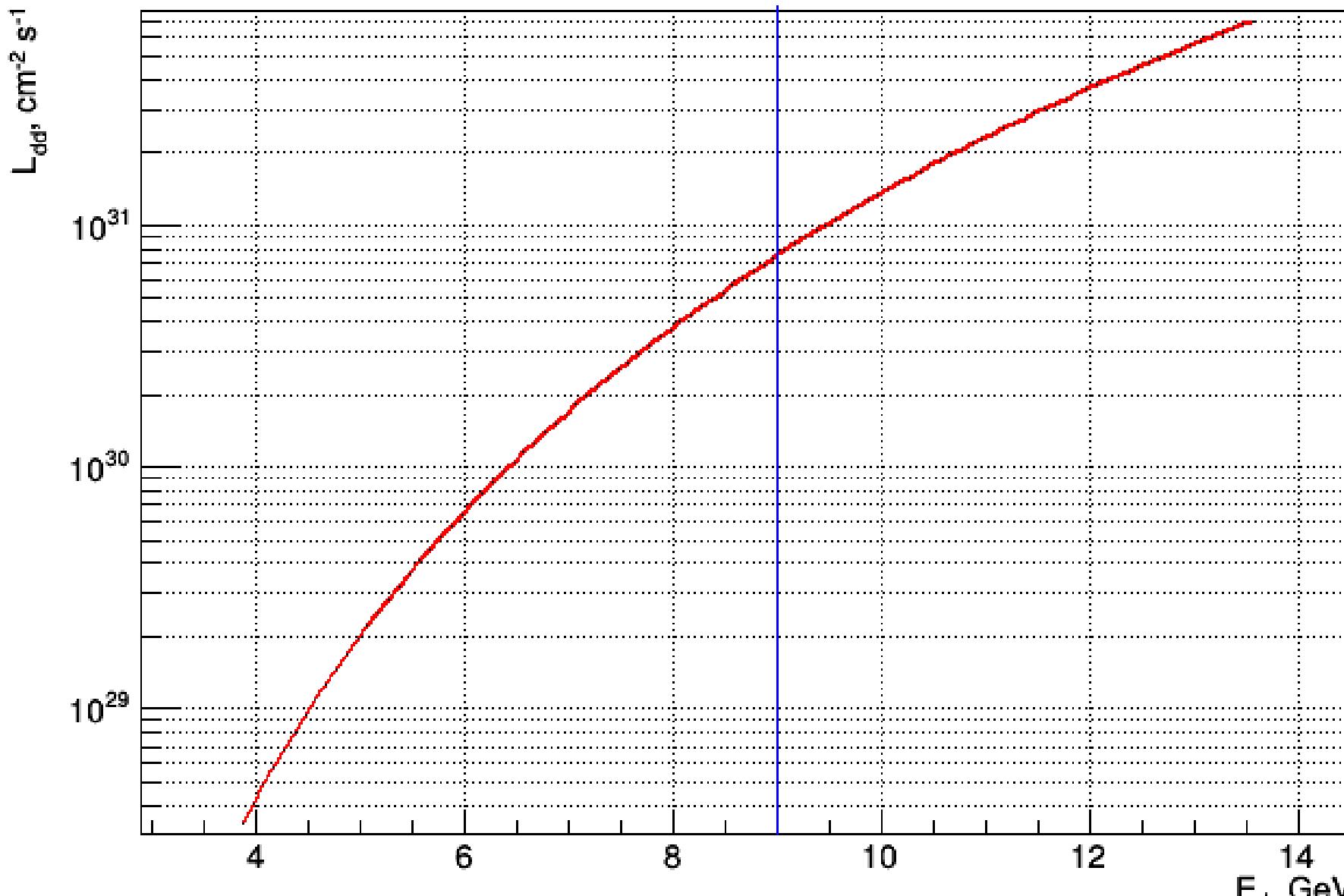
$$S(\Delta / 2) = \int_0^\infty [u^2(r) + w^2(r)] j_0(\Delta r / 2) r^2 dr$$

Deuteron elastic form factor



Inclusion of the deuteron form factor diminishes the cross section by ~ 2 orders of magnitude

Luminosity at SPD NICA for dd-collisions



Minimal energy of the deuteron in c.m.s. = 4GeV

Deuteron elastic form-factor is included

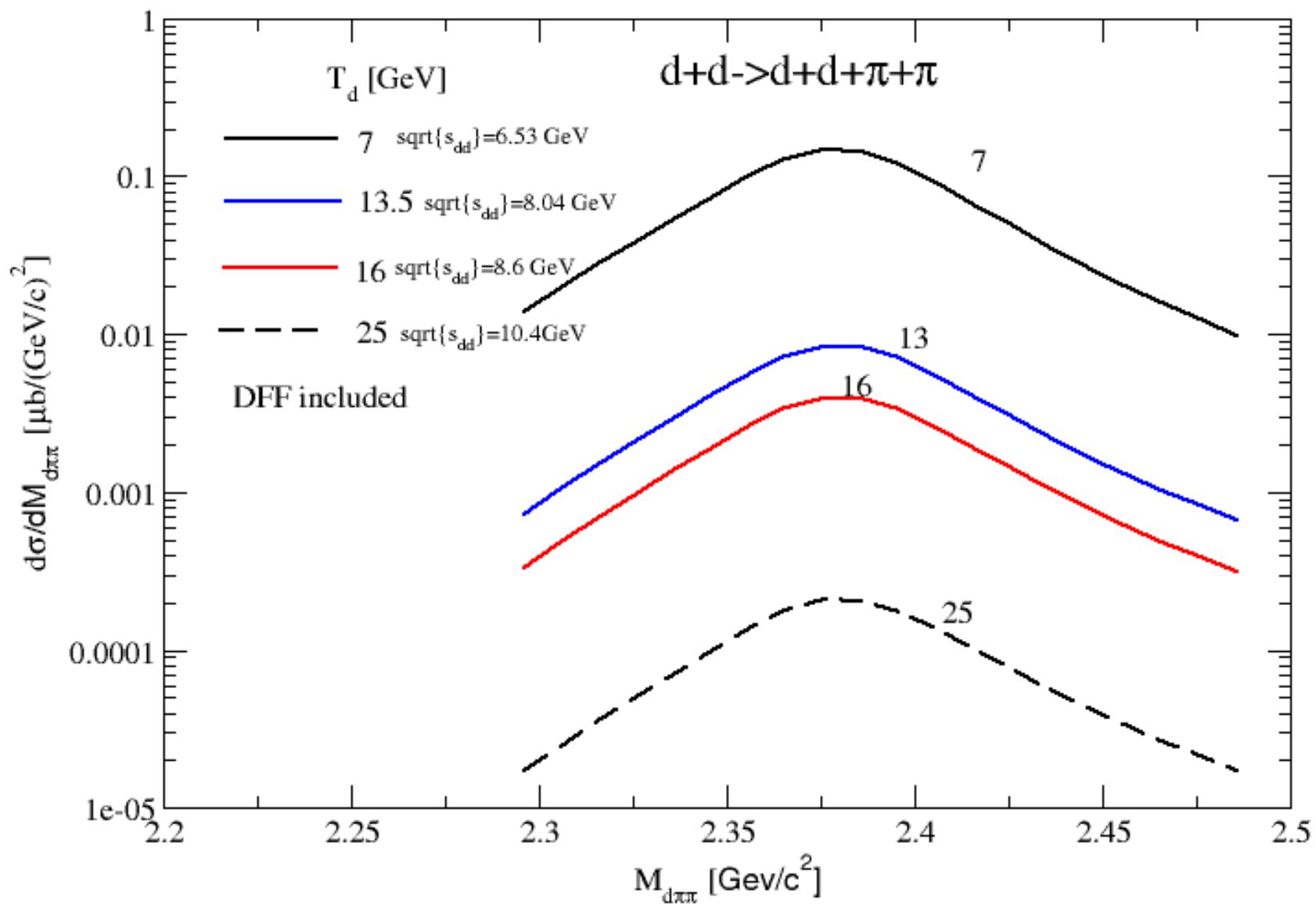
N- counting rate at maximum cross. sec.

$$E_d = 4 \text{ GeV}, L_{dd} = 4.1 \cdot 10^{28} \text{ cm}^{-2} \text{ s}^{-1}$$

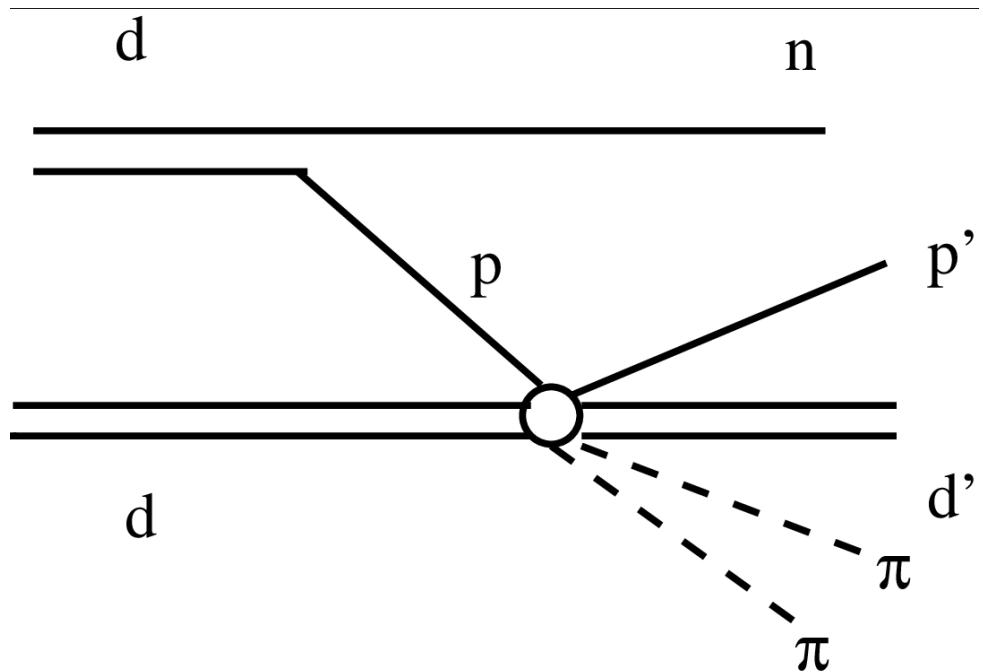
$$N = 0.14/\text{h}$$

$$E_d = 5.2 \text{ GeV}, L_{dd} = 2.1 \cdot 10^{29} \text{ cm}^{-2} \text{ s}^{-1}$$

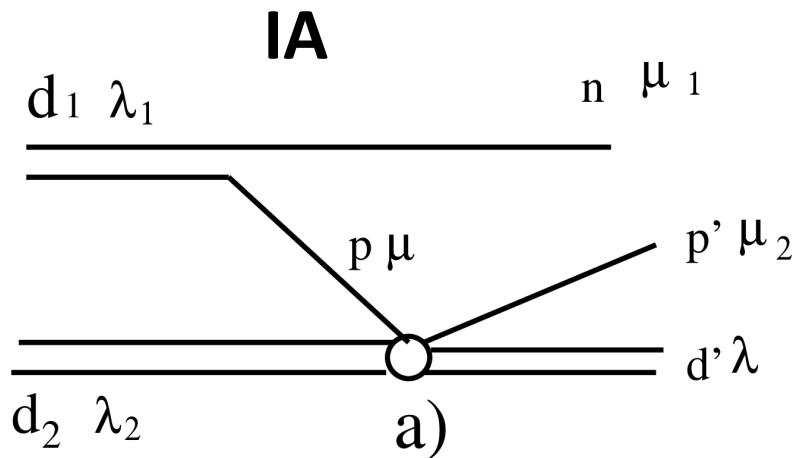
$$N = 0.007/\text{h}$$



Another variant: The $pd \rightarrow pd$ $\pi\pi$ subprocess in the
 $dd \rightarrow npd$ $\pi\pi$ reaction



Relations between $dd \rightarrow npd$ and $pd \rightarrow pd$ in the IA



S-wave dominates at $q < 0.15 \text{ GeV}/c$ and rescatterings are suppressed

Asymmetric pd-mode will be not available at NICA SPD, while symmetric dd mode will be established.

$$|M(dd \rightarrow npd)|^2 = K[u^2(q) + w^2(q)] |M(pd \rightarrow pd)|^2$$

d_2^\uparrow : Vector or tensor Polarized

$$A_Y^d(dd_2^\uparrow \rightarrow npd) = A_Y^d(pd^\uparrow \rightarrow pd),$$

$$A_{YY} = (dd_2^\uparrow \rightarrow npd) = A_{YY}(pd^\uparrow \rightarrow pd)$$

d_1^\uparrow : Vector Polarized

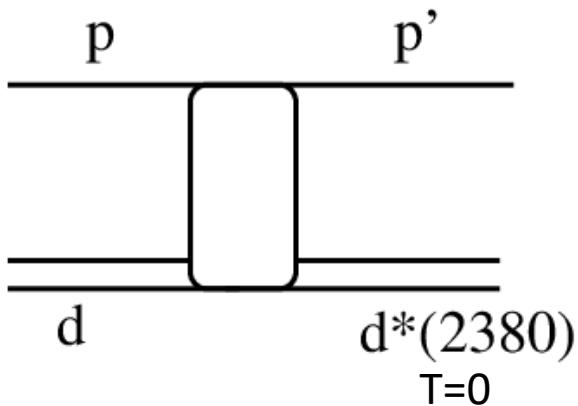
$$A_Y^d(d_1^\uparrow d \rightarrow npd) = \frac{2}{3} A_Y^p(p^\uparrow d \rightarrow pd)$$

Both d_1 and d_2 deuterons are vector or tensor polarized:

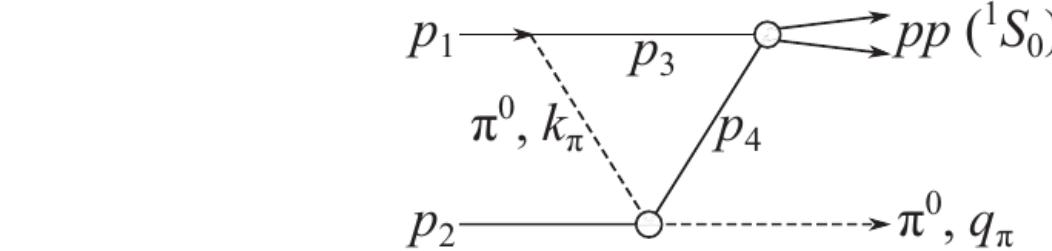
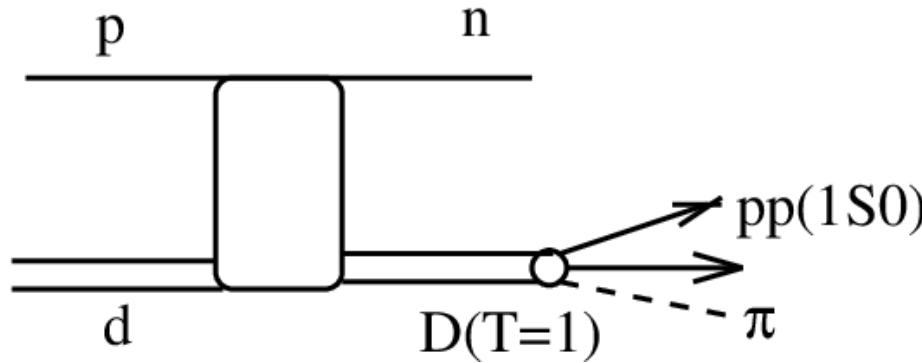
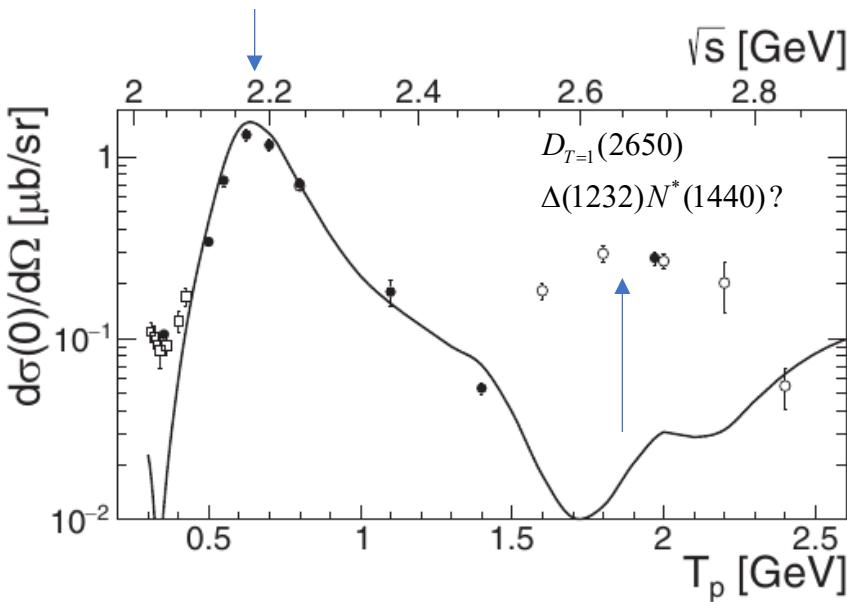
$$C_{Y,Y}(d^\uparrow d^\uparrow \rightarrow npd) = \frac{2}{3} C_{y,y}(p^\uparrow d^\uparrow \rightarrow pd)$$

$$C_{Y,YY}(d^\uparrow d^\uparrow \rightarrow npd) = \frac{1}{3} C_{y,yy}(p^\uparrow d^\uparrow \rightarrow pd)$$

ISOVECTOR DIBARYON RESONANCES: to search at SPD



PHYSICAL REVIEW C **107**, 015202 (2023)



PHYSICAL REVIEW C **107**, 015202 (2023)

Resonant behavior of the $pp \rightarrow \{pp\}_s \pi^0$ reaction at the energy $\sqrt{s} = 2.65$ GeV

D. Tsirkov^{1,2}, B. Baimurzinova^{1,2,3,*}, V. Komarov¹, A. Kulikov¹, A. Kunsafina^{1,2,3}, V. Kurbatov¹, Zh. Kurmanalyiev^{1,2,3} and Yu. Uzikov^{1,4,5}

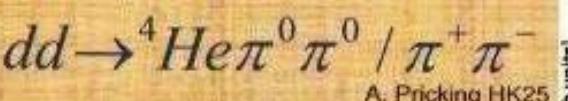
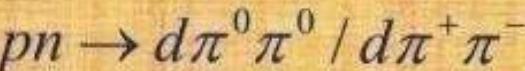
SUMMARY

- Properties of non-strange dibaryons are important for understanding (short-range) NN- dynamics. Compact objects or quasi-molecules?
- Isoscalar dibaryons d^* with $M=2380$ and $2470, 2630$ MeV can be searched at the first phase of the SPD NICA, in dd-mode at $\sqrt{s_{dd}} = 4-5$ GeV.
- The cross section of $dd \rightarrow d^*d \rightarrow dd \pi\pi$ is strongly suppressed by the DFF.
- The $pd \rightarrow pd\pi\pi$ is more preferable, but has to be extracted from the $dd \rightarrow n + pd\pi\pi$ (in the region of dominance of the IA)

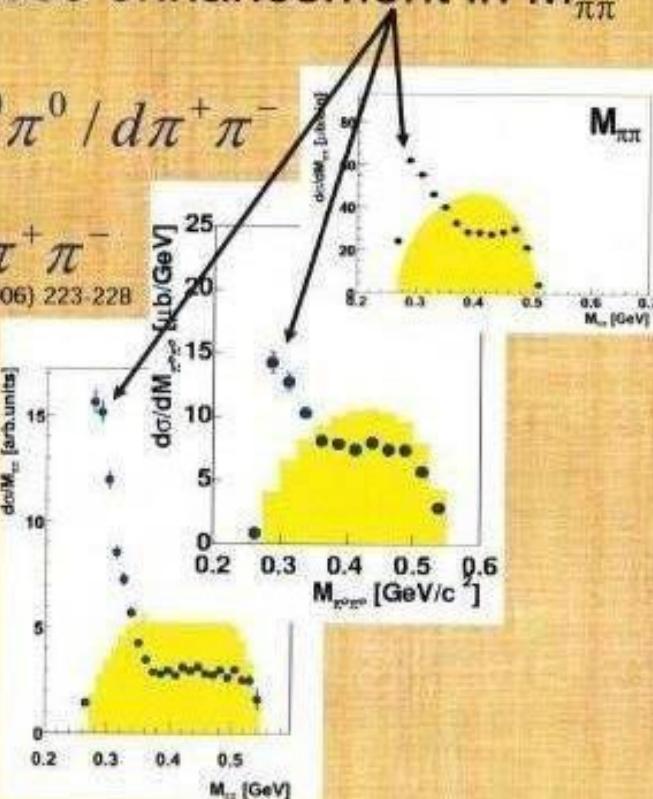
THANK YOU FOR ATTENTION!

History of the ABC effect

Low mass enhancement in $M_{\pi\pi}$

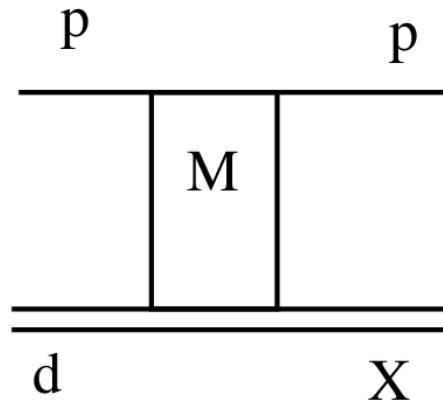


- Alexander Abashian,
Norman E. Booth
Kenneth M. Crowe,
Phys. Rev. Lett. **5**, 258 (1960)

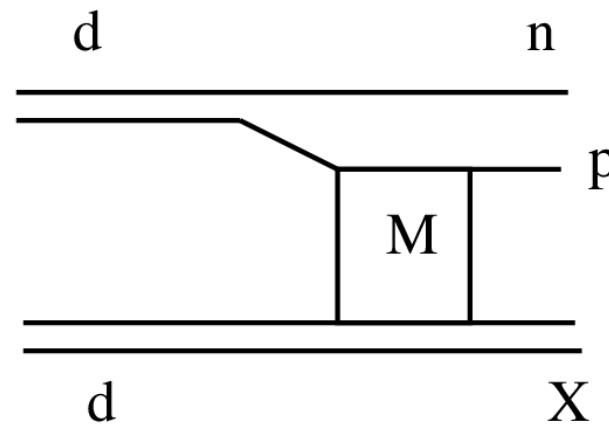


from talk by M. Bashkanov, DFG -meeting (Bonn,2010); $I_{\pi\pi} = 0 \implies \text{ABC}$
<http://wasasrv.ikp.kfa-juelich.de/WasaWiki/index.php/File:MB100716-ABC-meson10.pdf>

in Impulse approximation



a)



b)

$$T(dd \rightarrow n + pX) = \sum_{\sigma'} \langle \sigma_n, \sigma_{p'} | \psi_d^\lambda(\vec{q}) \rangle T_{\lambda\sigma'}^{M_X\sigma_p} (pd \rightarrow pX)$$

When the final neutron takes one half λ -wave dominance, suppressed p_T mode pd->pd amplitude can be extracted

$$\vec{p}_n = \vec{p}_d / 2$$

TABLE I. Parameters of resonances R and their decay channels $R \rightarrow a + b$ relevant for the $pp \rightarrow d(\pi\pi)_0$ reactions. For the parameter p_0 , the given interval corresponds to all possible isospin channels.

R	M_R (MeV)	$\Gamma_R^{(0)}$ (MeV)	ab	l	p_0 (MeV)	$\Gamma_{R \rightarrow ab}^{(0)}$ (MeV)	Λ_{ab} (GeV)
D_{03}	2376	77	np	2	730	9	0.35
			σd	2	350	2	0.18
			πD_{12}	1	173–176	31	0.12
D_{12}	2150	110	πd	1	221–223	33	0.15
Δ	1232	117	πN	1	226–229	117	0.16
σ	303	126	$\pi\pi$	0	72–80	126	0.09

M.P. Platonova, V.I. Kukulin,*Phys.Rev.D* 103 (2021) 11,
[1140250, Isospin symmetry breaking in double-pion production
 in the region of \$d^*\(2380\)\$ and the scalar \$\sigma\$ meson](#)

$$\Gamma_{R \rightarrow ab}(p) = \Gamma_{R \rightarrow ab}^{(0)} \left(\frac{p}{p_0} \right)^{2l+1} \left(\frac{p_0^2 + \Lambda_{ab}^2}{p^2 + \Lambda_{ab}^2} \right)^{l+1},$$

Resonant behavior of the $pp \rightarrow \{pp\}_s \pi^0$ reaction at the energy $\sqrt{s} = 2.65$ GeV

D. Tsirkov^{1,2}, B. Baimurzinova^{1,2,3,*}, V. Komarov,¹ A. Kulikov¹, A. Kunsafina^{1,2,3}, V. Kurbatov,¹ Zh. Kurmanalyiev^{1,2,3} and Yu. Uzikov^{1,4,5}

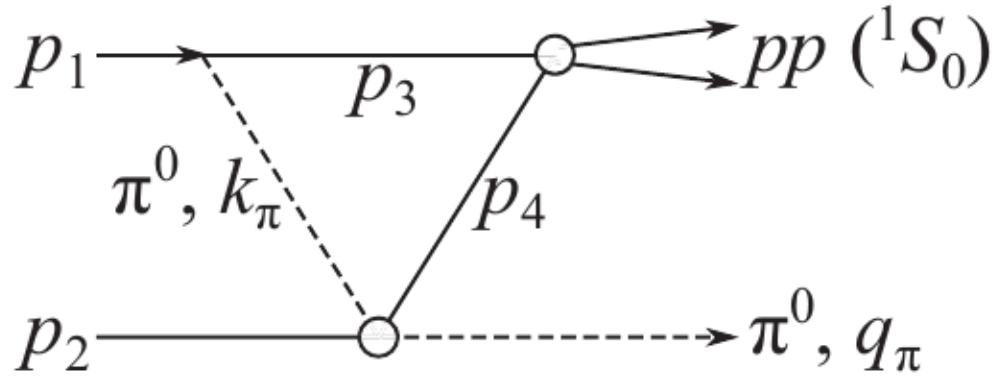


FIG. 6. The OPE mechanism of the reaction $pp \rightarrow \{pp\}_s \pi^0$.

shown. Note that the slope changes its sign in the region of the observed peak.

V. DISCUSSION

